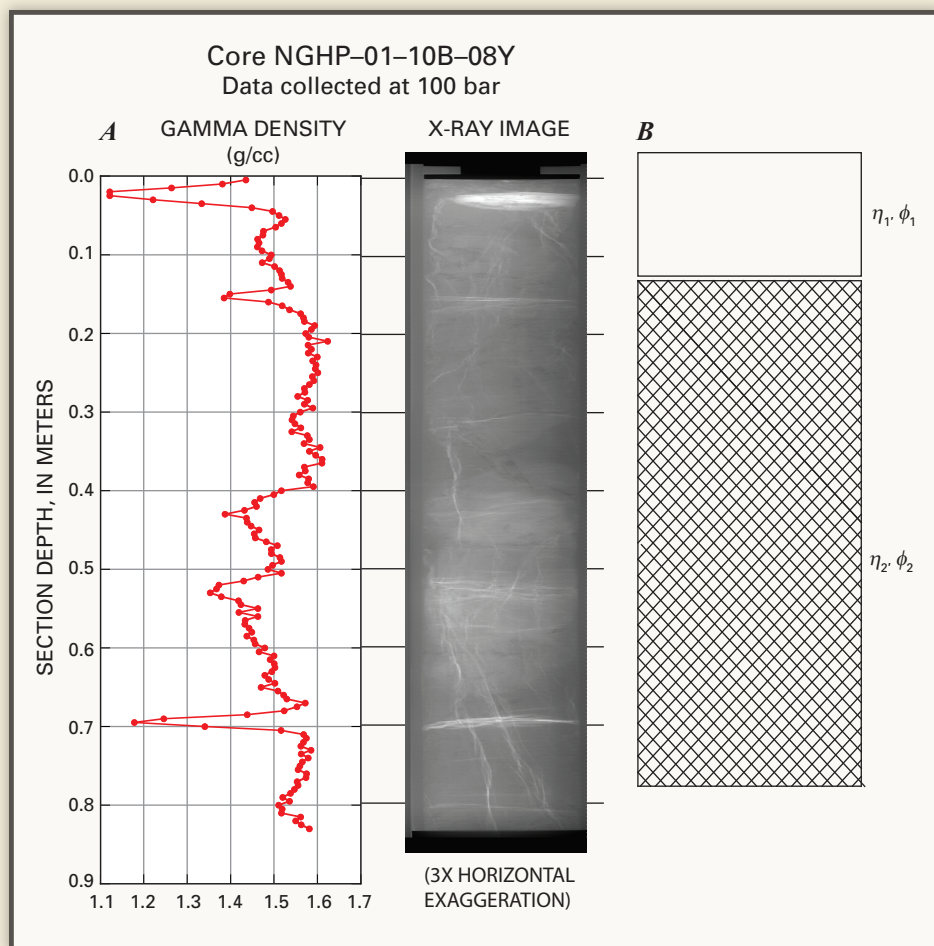


Anisotropic Velocities of Gas Hydrate-Bearing Sediments in Fractured Reservoirs



Scientific Investigations Report 2009–5141

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By Myung W. Lee

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Anisotropic Velocities of Gas Hydrate-Bearing Sediments in Fractured Reservoirs

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Abstract

During the Indian National Gas Hydrate Program Expedition 01 (NGHP-01), one of the richest marine gas hydrate accumulations was discovered at drill site NGHP-01-10 in the Krishna-Godavari Basin, offshore of southeast India. The occurrence of concentrated gas hydrate at this site is primarily controlled by the presence of fractures. Gas hydrate saturations estimated from P- and S-wave velocities, assuming that gas hydrate-bearing sediments (GHBS) are isotropic, are much higher than those estimated from the pressure cores. To reconcile this difference, an anisotropic GHBS model is developed and applied to estimate gas hydrate saturations. Gas hydrate saturations estimated from the P-wave velocities, assuming high-angle fractures, agree well with saturations estimated from the cores. An anisotropic GHBS model assuming two-component laminated media—one component is fracture filled with 100-percent gas hydrate, and the other component is the isotropic water-saturated sediment—adequately predicts anisotropic velocities at the research site.

Introduction

Gas hydrates are a type of clathrate in which water molecules surround gas molecules to form an ice-like crystalline solid. In situ physical characteristics of gas hydrate-bearing sediments (GHBS) have been investigated extensively because of their widespread occurrence in most of the world oceans, as well as in permafrost regions (Kvenvolden, 1993), and their recognition as a negative feedback control on global temperature fluctuations (Archer, 2007), as a potential energy resource (Ruppel, 2007), and as a factor in sea-floor stability and safety issues (Nixon and Grozic, 2007).

As part of the Indian energy resource program, the Indian National Gas Hydrate Program Expedition 01 (NGHP-01) in 2006 was designed to investigate the geologic and geochemical controls on gas hydrate occurrences offshore of the Indian Peninsula and along the Andaman convergent margin.

A general description on geological and geophysical studies related to the gas hydrate occurrence in the offshore regions of India is well documented in Rao (2001). A total of 39 holes were drilled at 21 sites (fig. 1), some of which were cored as part of the expedition. One of the richest marine gas hydrate accumulations was discovered at drill site NGHP-01-10 in the Krishna-Godavari Basin (Collett and others, 2008). The initial site report used logging data and infrared images of recovered cores to infer the presence of gas hydrate within the depth interval from ≈ 25 to ≈ 160 m below the sea floor (mbsf). The gas hydrate existed as solid nodules, as fracture fill in high-angle and subhorizontal veins, and was disseminated throughout the cores.

The gas hydrate saturations estimated from the logging-while-drilling (LWD) resistivity log reached as much as 80 percent of the pore space (Collett and others, 2008), whereas the highest saturations estimated from the pressure core at the nearby wells are about 26 percent of the pore space. The large difference of gas hydrate estimations is speculated to be caused by the anisotropic nature of the reservoir due to fractures at this site (Lee and Collett, 2009).

Well logs have been extensively used to characterize the in situ gas hydrate-bearing sediments including their saturations (for example, Collett, 2000, 2002; Guerin and others, 1999; Helgerud and others, 1999; Hyndman and others, 2001; Lee and Collett, 2005). The majority of investigations have been focused on the unfractured sandstone reservoirs and dispersed gas hydrate in the shale, but few investigations analyzed well logs acquired within the fractured reservoirs.

The purpose of this report is to derive a realistic fracture model to predict P- and S-wave velocities and to compare gas hydrate saturations estimated from the P- and S-wave velocities with those estimated from pressure cores by incorporating anisotropic properties of fractured reservoir. Transverse isotropic theory for laminated media (for example, Backus, 1962; Postma, 1955; White and Angona, 1955) and anisotropy theory by Thomsen (1986, 1995) are used to describe elastic velocities for the fractured reservoir, and gas hydrate saturations are estimated from P- and S-wave velocities assuming isotropic and anisotropic GHBS.

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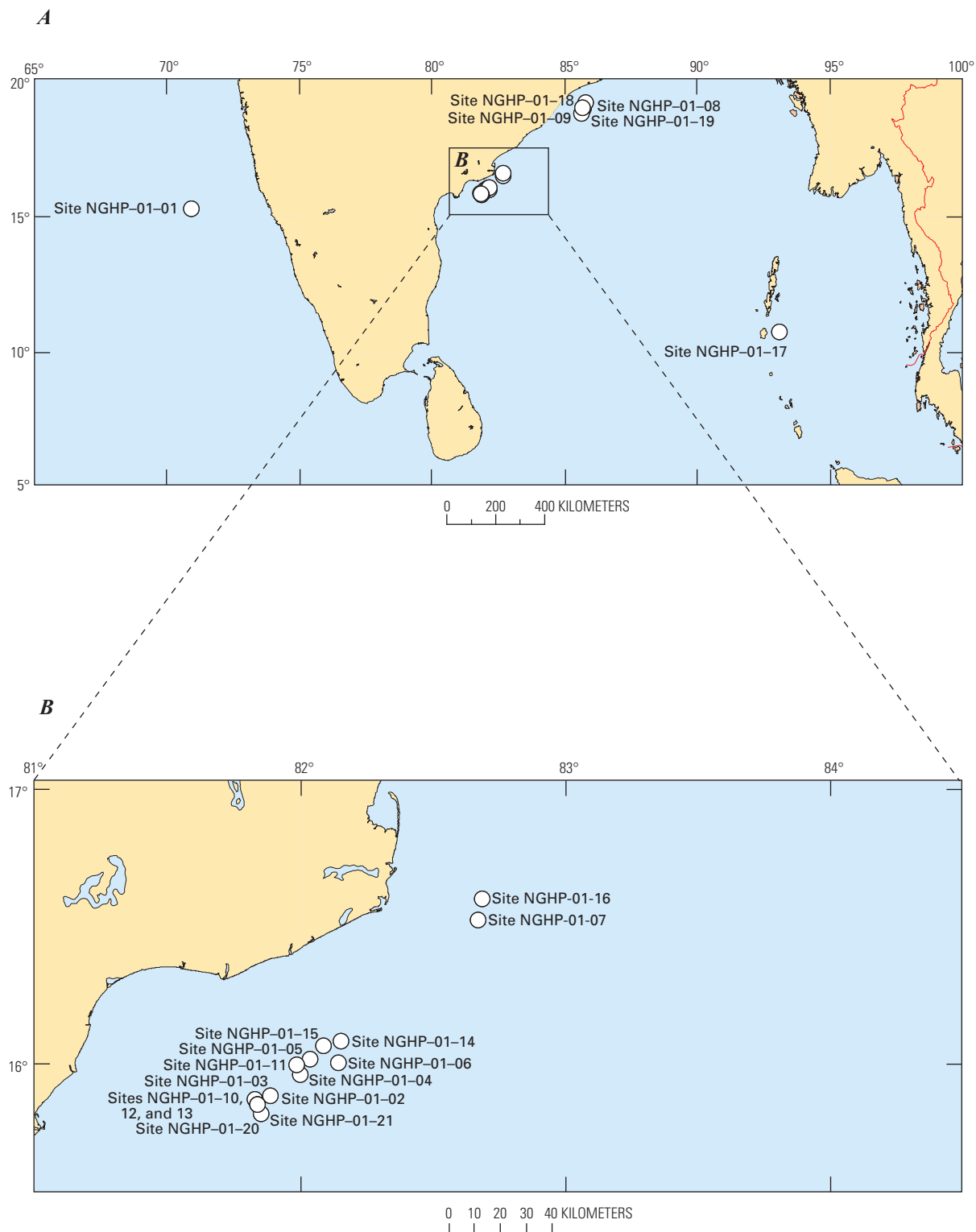


Figure 1. A, NGHP Expedition 01 site map depicting the location of the research drill sites established during the expedition. B, Insert map of the NGHP Expedition 01 drill sites in the Krishna-Godavari Basin. Modified from Collett and others (2008).

Theory

Isotropic Velocities

Velocities of isotropic GHBS, whereby the gas hydrate fills the sediment's pore space, are calculated using the three-phase Biot equation (Lee, 2007). Lee (2008) derived the following bulk (k) and shear moduli (μ) of the GHBS applicable for low-frequency well-log and seismic data:

$$\begin{aligned} k &= K_{ma}(1 - \beta_p) + \beta_p^2 K_{av} \\ \mu &= \mu_{ma}(1 - \beta_s) \end{aligned} \quad (1)$$

with

$$\frac{1}{K_{av}} = \frac{(\beta_p - \phi)}{K_{ma}} + \frac{\phi_w}{K_w} + \frac{\phi_h}{K_h},$$

and

$$\begin{aligned} \beta_p &= \frac{\phi_{as}(1 + \alpha)}{(1 + \alpha\phi_{as})}, \\ \beta_s &= \frac{\phi_{as}(1 + \gamma\alpha)}{(1 + \gamma\alpha\phi_{as})} \end{aligned}$$

where α is the consolidation parameter (Pride and others, 2004; Lee, 2005) with $\gamma = \frac{1 + 2\alpha}{1 + \alpha}$; $\phi_{as} = \phi_w + \varepsilon\phi_h$, $\phi_w = (1 - C_h)\phi$, and $\phi_h = C_h\phi$. K_{ma} , K_w , and K_h are the bulk moduli of grain, water, and gas hydrate respectively; μ_{ma} is shear modulus of the grain.

The parameter ε accounts for the reduced impact of hydrate formation relative to compaction in terms of stiffening the host-sediment framework, and Lee and Waite (2008) recommended $\varepsilon = 0.12$ for modeling velocities of GHBS.

The P-wave (V_p) and S-wave velocities (V_s) of the GHBS are given by:

$$V_p = \sqrt{\frac{k + 4\mu/3}{\rho_b}} \text{ and } V_s = \sqrt{\frac{\mu}{\rho_b}} \quad (2)$$

where

$$\begin{aligned} \rho_b &\text{ is the bulk density of GHBS given by} \\ \rho_b &= \rho_s(1 - \phi) + \rho_w\phi(1 - C_h) + \rho_h\phi C_h. \end{aligned}$$

The use of equations 1 and 2 for computing velocities of GHBS is referred to as the simplified three-phase equation (STPE) in this report. For water-saturated sediments, STPE is the same as the Gassmann equation.

The consolidation parameter α depends on the effective pressure and degree of consolidation. Mindlin (1949) showed that the bulk and shear moduli depend on the one-third power of the effective pressure. On the basis of the theory by Mindlin (1949), the depth (or effective pressure) dependent α is proposed by the following equation using the exponent n :

$$\alpha_i = \alpha_0(p_0 / p_i)^n \approx \alpha_0(d_0 / d_i)^n \quad (3)$$

where

α_0 is the consolidation parameter at the effective pressure p_0 or depth d_0 ,

and

α_i is the consolidation parameter at the effective pressure p_i or depth d_i .

The use of $n = 1/3$ in equation (3) is similar to the Mindlin relation.

Anisotropic Velocities

Out of 21 holes drilled in the Indian Ocean during NGHP-01, Collett and others (2008) identified four gas hydrate accumulations within the fractured shale. Figure 2A shows an X-ray image for recovered core at drill site NGHP-01-10B at the depth of 50.1 mbsf. The light color in the image denotes the presence of gas hydrate, and figure 2A shows five or six layers of gas hydrate and wispy subvertical veins within a relatively homogeneous matrix (Collett and others, 2008). Figure 2B shows a schematic diagram for modeling the fractured reservoir.

Elastic anisotropy due to aligned cracks in porous rock has been investigated extensively (for example, Eshelby, 1957; Walsh, 1969; Watt and others, 1976; Hudson, 1980; Thomsen, 1995). Because it is assumed that the fracture is filled with 100-percent gas hydrate, the fractured reservoir can be modeled as a laminated elastic media (White and Angona, 1955) rather than saturated aligned cracks embedded in porous media (for example, see Thomsen, 1995). In this report, a fractured reservoir is modeled as a layered media composed of two components as shown in figure 2B. The first component is fractures filled with 100-percent gas hydrate, and its characteristics are given by P-wave velocity (V_{p1}), S-wave velocity (V_{s1}), and density (ρ_1); the second component is isotropic medium with P-wave velocity (V_{p2}), S-wave velocity (V_{s2}), and density (ρ_2). Anisotropic P- and S-wave velocities can be computed using equations shown in White (1965). Because the fracture is filled with 100-percent gas hydrate, the elastic constants for the fracture are identical to those of the pure gas hydrate.

Let's define η_1 and η_2 to be the volume fraction of component 1 and volume fraction of component 2, respectively. Phase velocities of transverse isotropic media caused from fractures can be computed using the following definition:

$$\langle G \rangle \equiv \eta_1 G_1 + \eta_2 G_2$$

$$\langle G \rangle^{-1} \equiv \left(\frac{\eta_1}{G_1} + \frac{\eta_2}{G_2} \right)^{-1}$$

where

G is any elastic constant such as density of the component 1 and component 2.

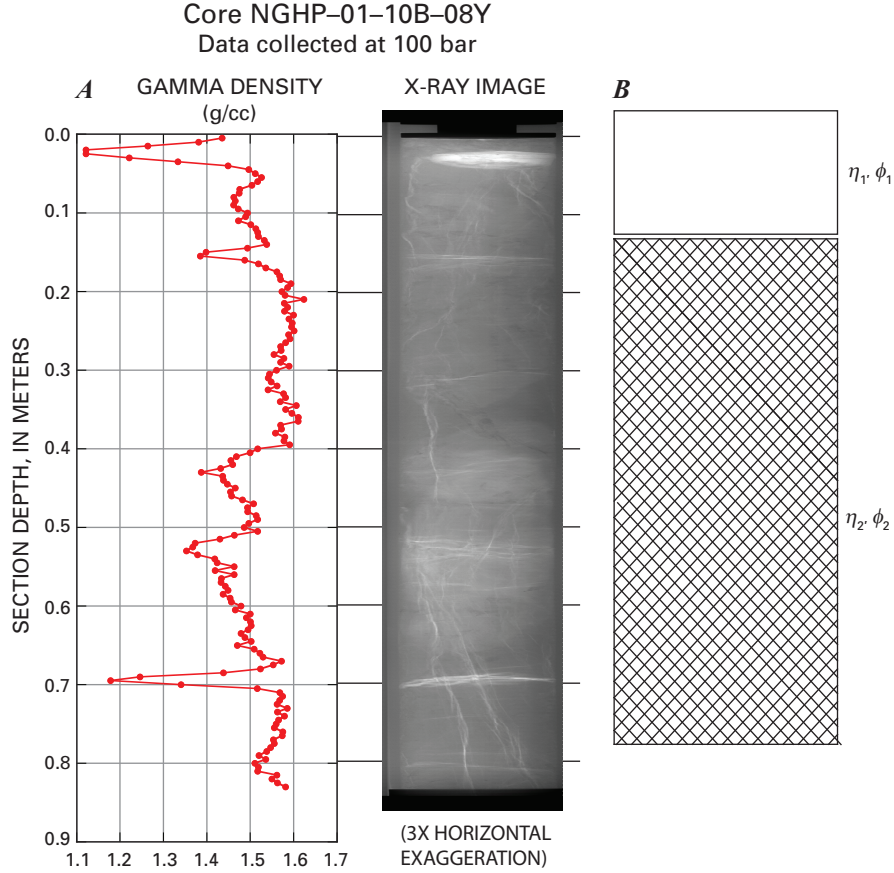


Figure 2. A, X-ray image of gas hydrate-bearing sediments at the depth of 50.1 meters below sea floor at the drill site NGHP-01-10B (modified from Collet and others, 2008). B, Schematic diagram showing laminated two-component sediment package. Parameter η_1 is the volume fraction of component 1, η_2 is the volume fraction of component 2, and ϕ is porosity.

The P- and S-wave velocities of transverse isotropic media can be calculated from the following equations using the Lamé constants λ and μ (White, 1965):

$$A = \left\langle \frac{4\mu(\lambda + \mu)}{(\lambda + 2\mu)} \right\rangle + \left\langle \frac{1}{(\lambda + 2\mu)} \right\rangle^{-1} \left\langle \frac{y}{(\lambda + 2\mu)} \right\rangle^2 \quad (4)$$

$$C = \left\langle \frac{1}{(\lambda + 2\mu)} \right\rangle^{-1}$$

$$F = \left\langle \frac{1}{(\lambda + 2\mu)} \right\rangle^{-1} \left\langle \frac{\lambda}{(\lambda + 2\mu)} \right\rangle$$

$$L = \left\langle \frac{1}{\mu} \right\rangle^{-1}$$

$$N = \langle \mu \rangle$$

$$\rho = \langle \rho \rangle$$

$$V_p = \left(\frac{A \sin^2 \varphi + C \cos^2 \varphi + L + Q}{2\rho} \right)^{1/2}$$

$$V_s^V = \left(\frac{A \sin^2 \varphi + C \cos^2 \varphi + L + Q}{2\rho} \right)^{1/2}$$

$$V_s^H = \left(\frac{N \sin^2 \varphi + L \cos^2 \varphi}{\rho} \right)^{1/2}$$

$$Q = \sqrt{[(A - L) \sin^2 \varphi - (C - L) \cos^2 \varphi]^2 + 4(F + L)^2 \sin^2 \varphi \cos^2 \varphi}$$

where

- φ is the angle between the wavefront normal and the vertical axis (perpendicular to the layering);
- V_p is the anisotropic P-wave velocity,
- V_s^H is the horizontally polarized S-wave (SH) velocity,
- and
- V_s^V is the vertically polarized S-wave (SV) velocity.

Generally, transverse isotropic elastic media is described using Thomsen's parameters. Thomsen (1986, 1995) described an anisotropic elastic media using the three parameters γ , δ , and ϵ . Using these parameters, velocities are given by:

$$\begin{aligned}
V_p^2 &= \alpha_o^2(1 + 2\delta \sin^2 \varphi + 2\varepsilon \sin^4 \varphi) \\
(V_s^H)^2 &= \beta_o^2(1 + 2\gamma \sin^2 \varphi) \\
(V_s^V)^2 &= \beta_o^2 \left[1 + \frac{2\alpha_o^2(\varepsilon - \delta) \sin^2 \varphi \cos^2 \varphi}{\beta_o^2} \right]
\end{aligned} \tag{5}$$

where α_o and β_o are the P- and S-wave velocities at $\phi = 0$ or in the symmetry direction.

Thomsen's parameters are related to White's equation by the following equations (Thomsen, 1986):

$$\begin{aligned}
\gamma &= \frac{N - L}{2L} \\
\delta &= \frac{(F + L)^2 - (C - L)^2}{2C(C - L)} \\
\varepsilon &= \frac{A - C}{2C}
\end{aligned} \tag{6}$$

The velocities shown in equations 4 and 5 are phase velocities. Unless incident angle is 0° or 90° , the phase and group velocities are different. If the dip of the fracture is not 0° or 90° , the group velocities should be used to calculate the saturations.

Thomsen (1986) derived the relation between group and phase velocities for transversely isotropic media. Let φ_g be the angle between the ray direction and the vertical axis and GV_p , GV_s^H , and GV_s^V be the group velocities of P-wave, horizontally polarized shear-wave, and vertically polarized S-wave velocities, respectively. Then

$$\begin{aligned}
V_p(\varphi) &= GV_p(\varphi_g) \\
V_s^H(\varphi) &= GV_s^H(\varphi_g) \\
V_s^V(\varphi) &= GV_s^V(\varphi_g)
\end{aligned} \tag{7}$$

Note that equation 7 does not mean that group velocity equals phase velocity. Equation 7 indicates that if the corresponding wavefront normal angle φ is calculated using the following formula for a given ray angle φ_g , then equation 7 is applicable to compute group velocity:

$$\tan \varphi_g = \frac{\left(\tan \varphi + \frac{1}{V} \frac{dV}{d\varphi} \right)}{\left(1 - \frac{\tan \varphi}{V} \frac{dV}{d\varphi} \right)} \tag{8}$$

For weakly anisotropic media, group velocities are (Thomsen, 1986):

$$\begin{aligned}
\text{For P-wave: } \tan \varphi_g &= \tan \varphi [1 + 2\delta + 4(\varepsilon - \delta) \sin^2 \varphi] \\
\text{For SH-wave: } \tan \varphi_g &= \tan \varphi [1 + 2\gamma] \\
\text{For SV-wave: } \tan \varphi_g &= \tan \varphi \left[1 + \frac{2\alpha_o^2(\varepsilon - \delta)(1 - 2\sin^2 \varphi)}{\beta_o^2} \right]
\end{aligned} \tag{9}$$

For more details for the group and phase velocities, consult Thomsen (1986).

Modeling

Fractures in High-Porosity Shale

Figure 3 shows modeled P- and S-wave velocities for a fractured shale reservoir using the elastic constants shown in table 1. The model parameters for the fracture filled with 100-percent gas hydrate are $V_{p1} = 3.744$ km/s, $V_{s1} = 1.946$ km/s, and $\rho_1 = 0.926$ g/cm³; parameters for 100-percent water-saturated sediment are $V_{p2} = 1.503$ km/s, $V_{s2} = 0.166$ km/s, and $\rho_2 = 1.563$ g/cm³; $\phi_2 = 0.65$ and $C_v = 0.6$, where C_v is clay volume fraction. Velocities of water-saturated sediments were calculated using the STPE with $\alpha = 100$. Velocities of shale in this fracture model represent the velocities of shallow marine sediments.

Figure 3 shows a polar plot of P- and S-wave phase velocities with respect to the phase angle for a horizontal fracture whose volume fraction is 15 percent of the sediment volume, or the volume of a fracture filled with 100-percent gas hydrate is 15 percent of the total sediment volume. The P-wave velocity at the phase angle of 0° or equivalent for a horizontal fracture is 1.64 km/s, whereas the P-wave velocity at the phase angle of 90° or equivalent for vertical fractures is 1.83 km/s (about an 11-percent increase). The minimum SH-wave velocity of 0.186 km/s occurs at the phase angle of 0° , whereas the maximum SH-wave velocity of 0.621 km/s occurs at the phase angle of 90° (about a 200-percent increase). These anisotropic velocities can be calculated using Thomsen's parameters using equation 5, and these parameters are $\gamma = 5.09$, $\delta = -0.0729$, and $\varepsilon = 0.1169$, which are calculated from equation 6. The magnitude of these parameters increases with increasing fracture density (or gas hydrate volume).

For comparison, velocities for isotropic GHBS (gas hydrate accumulates in the pore space of the sediments) having 15-percent gas hydrate volume content are shown in figure 3. Velocities are computed using the STPE with a constant $\alpha = 100$ and the same parameters for the water-saturated sediments for the fracture model. Depending on the phase angle, the isotropic velocities are smaller or larger than the anisotropic velocities.

Figure 4 shows calculated velocities for isotropic GHBS along with transversely isotropic velocities for the fractured reservoir with respect to volume fraction of gas hydrate. Isotropic velocities of GHBS are calculated from the STPE with $\alpha = 100$ and $C_v = 0.6$, which are the same as those for water-saturated sediment of the fracture model. For isotropic medium, the volume fraction of gas hydrate (V_h) is given by $V_h = \phi S_h$, where S_h is the gas hydrate saturation in the pore space and is given by the volume fraction of fractures for a fractured reservoir. In the case that there is no gas hydrate either in the pore or fracture, the two velocities are the same. Because the maximum volume fraction for the isotropic medium is limited by porosity, maximum volume fraction for the isotropic GHBS is 0.65. For 100-percent gas hydrate saturation in the fracture, the velocities are velocities of pure methane gas hydrate.

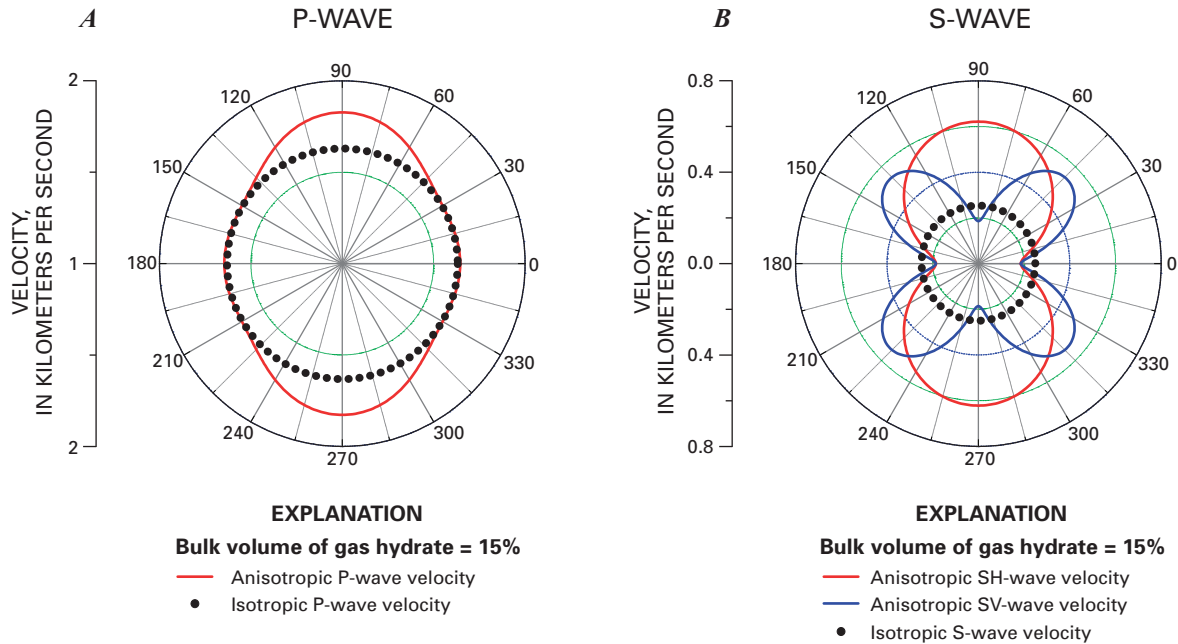


Figure 3. Isotropic velocities of gas hydrate-bearing sediment and anisotropic velocities for horizontal fracture having a 15-percent bulk volume of gas hydrate with respect to the angle of incident. The model represents fractures in high-porosity shale. *A*, P-wave velocity. *B*, Horizontally polarized (V_s^H) and vertically polarized S-wave velocity (V_s^V).

As shown in figure 4*A*, the effect of anisotropy is much more pronounced on the S-wave velocities as expected from figure 3. When the fracture orientation is perpendicular to the direction of measurement (incident angle is 0°), the S-wave velocities are almost constant except at very high volumes of gas hydrate saturations above 80 percent. On the other hand, in the case that the fracture is parallel to the direction of measurement, the S-wave velocities increase rapidly at low saturation and increase monotonically as gas hydrate saturation further increases. For all ranges of fracture volume, velocities modeled for isotropic GHBS lie between two velocities modeled for anisotropic GHBS.

Figure 4*B* shows the relation between modeled anisotropic P- and S-wave velocities along with the measured relation at drill site NGHP-01-10D. Based on the theoretical

curves for P- and S-wave velocities by angle of incidence, the measured data do not exhibit isotropic behavior. A series of theoretical curves were calculated to determine the fracture orientation representative of the measured data. Assuming a fracture dip of 85° , or ray angle of 85° , the modeled P- and S-wave velocity relation is similar to the measured data, especially within the shaded area.

The velocity at the ray angle of 85° is the group velocity, not the phase velocity shown in figure 3. Because energy is transmitted with the group velocity, the group velocity should be used to analyze measured velocities. As shown in equation 9, the relation between the ray angle and phase angle depends on Thomsen's parameters, which depend on the fracture volume and elastic properties of sediments encasing the fracture. Consequently, for a given ray angle, the phase angle to be used in calculating the group velocity varies with gas hydrate saturation. Assuming the fracture volume is 15 percent of the sediment (same as for figure 3), group and phase velocities using equations 7 and 8 are calculated and are shown in figure 5.

Because the P-wave velocity has only a small amount of anisotropy, the phase and group velocities are similar to those shown in figure 5*A*. In other words, the phase angle can be interchanged into the ray angle without significant error in the P-wave velocity. However, group and phase velocities for the horizontally polarized shear wave are quite different as shown in figure 5*B*. For ray angles less than about 60° , the group velocities are almost the same. A greater change of group velocity occurs for ray angles greater than 60° . In other words, the SH-wave velocities are insensitive to the

Table 1. Constants used for modeling.

[Subscript *s*, *h*, *c*, and *w* are sand, gas hydrate, clay, and water, respectively; *K*, μ , and ρ are bulk modulus, shear modulus, and density, respectively. The bulk and shear moduli of gas hydrate are calculated at 11°C and 1 MPa from Helgerud and others (2009) and others are from Lee (2005). The bulk and shear moduli of gas hydrate yield the P- and S-wave velocities of pure gas hydrate as 3.744 km/s and 1.946 km/s, respectively. MPa; Mega Pascals, GPa, Giga Pascals]

Bulk modulus	Shear modulus	Density
$K_s = 38 \text{ GPa}$	$\mu_s = 44 \text{ GPa}$	$\rho_s = 2,650 \text{ kg/m}^3$
$K_c = 20.9 \text{ GPa}$	$\mu_c = 6.85 \text{ GPa}$	$\rho_c = 2,580 \text{ kg/m}^3$
$K_h = 8.27 \text{ GPa}$	$\mu_h = 3.49 \text{ GPa}$	$\rho_h = 922 \text{ kg/m}^3$
$K_w = 2.29 \text{ GPa}$	$\mu_w = 0$	$\rho_w = 1,000 \text{ kg/m}^3$

dip of the fracture, if the dip is less than $\approx 60^\circ$. However, this observation for SH-waves cannot be generalized, because the relation between ray angle and SH-velocity is a complicated function of porosity and fracture density.

Fractures in Sandstone

Figure 6 shows modeled P- and S-wave velocities for a fractured reservoir in sandstone using elastic constants shown in table 1. The model parameters for the fracture filled with 100-percent gas hydrate in sandstone are identical to those for the fracture in shale. Parameters for 100-percent water-saturated sandstone are $V_{p2} = 1.950$ km/s, $V_{s2} = 0.650$ km/s, $\rho_2 = 1.986$ g/cm³, $\phi_2 = 0.4$, and $C_{v2} = 0.1$. Velocities of sandstone in this fracture model represent the velocities of host sediments at depths of around 600–800 m, typical host sediments of GHBS at the Alaska North Slope (Lee, 2005).

Figure 6 show a polar plot of P- and S-wave velocities with respect to the angle of incidence for a horizontal fracture whose volume fraction is 15 percent. The P-wave velocity at the incident angle of 0° is 2.108 km/s, whereas the P-wave

velocity at the incident angle of 90° is 2.144 km/s (less than 2-percent increase). The SH-wave velocity at the incident angle of 0° is 0.720 km/s, whereas the SH-wave velocity at the incident angle of 90° is 0.826 km/s (about a 16-percent increase). These anisotropic velocities can be calculated using Thomsen's parameters using equation 5, and these parameters are calculated using equation 6 as $\gamma = 0.01573$, $\delta = -0.0340$, and $\varepsilon = 0.0204$.

For comparison, velocities for isotropic GHBS having 15-percent gas hydrate volume content are shown in figure 6. Velocities of isotropic GHBS are computed using the STPE with $\alpha = 30$ and the same parameters for the water-saturated sediments used for the fracture model (that is $\rho_2 = 1.986$ g/cm³, $\phi_2 = 0.4$, and $C_{v2} = 0.1$). The isotropic P-wave velocity is 2.347 km/s and S-wave velocity is 0.871 km/s, which are larger than the maximum anisotropic P- and S-wave velocities. Unlike velocities for the fractured shale reservoir, velocities of anisotropic velocities in sandstone are insensitive to the orientation of fractures.

Figure 7A shows calculated velocities for isotropic GHBS along with transversely isotropic velocities for the fracture reservoir in sandstone with respect to volume fraction

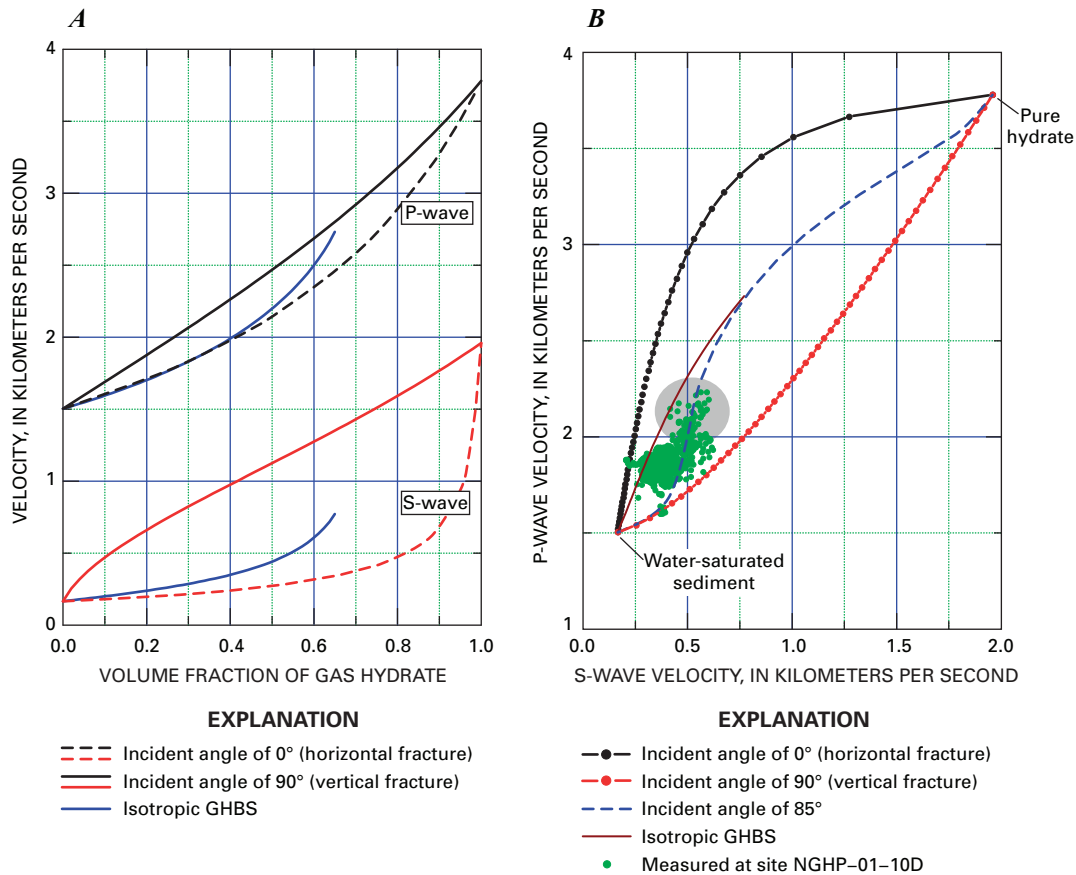


Figure 4. Modeling elastic velocities for anisotropic gas hydrate-bearing sediments (GHBS) caused by fracture in high-porosity shale and for isotropic GHBS. A, Relation between gas hydrate saturation and bulk volume of gas hydrate in sediments. B, Relation between S-wave velocity and P-wave velocity.

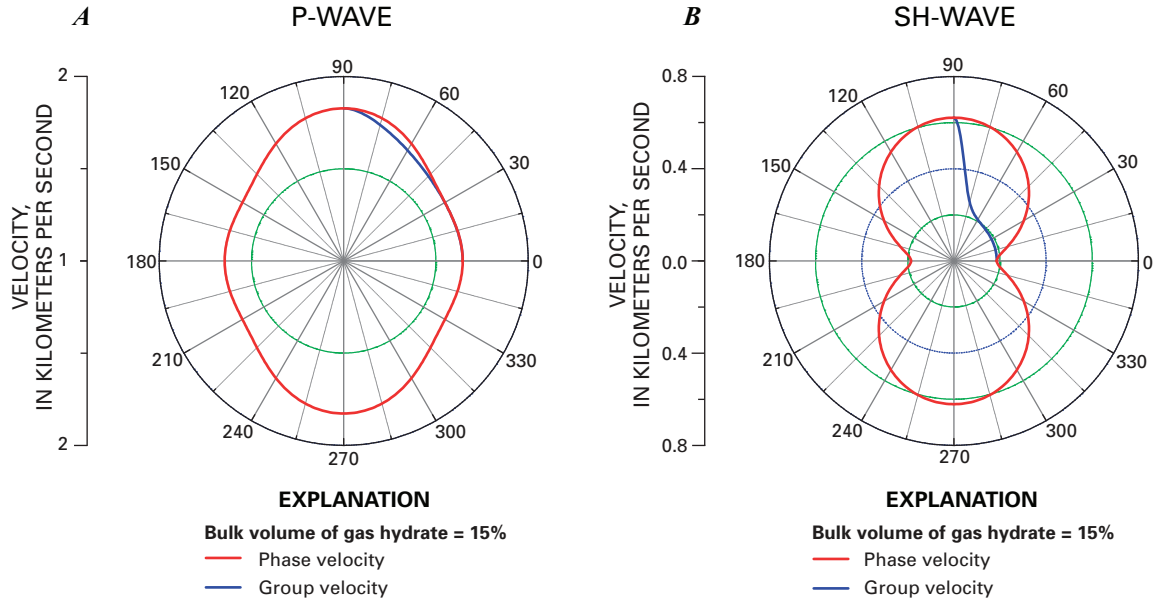


Figure 5. Group and phase velocities for a horizontal fracture having 15-percent bulk volume of gas hydrate (or fracture volume) with respect to phase and ray angles. *A*, P-wave velocity. *B*, SH-wave velocity.

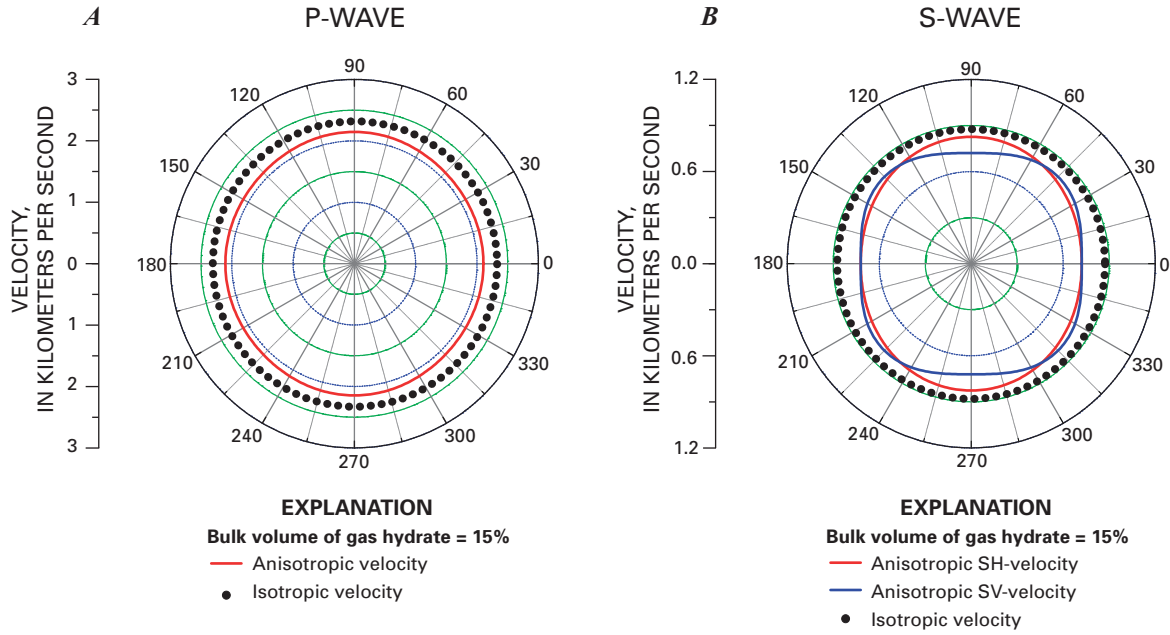


Figure 6. Isotropic velocities of gas hydrate-bearing sediment and anisotropic velocities for horizontal fracture having a 15-percent bulk volume of gas hydrate with respect to the angle of incidence. The model represents fractures in clean sandstone. *A*, P-wave velocity. *B*, Horizontally polarized (V_s^h) and vertically polarized S-wave velocity (V_s^v).

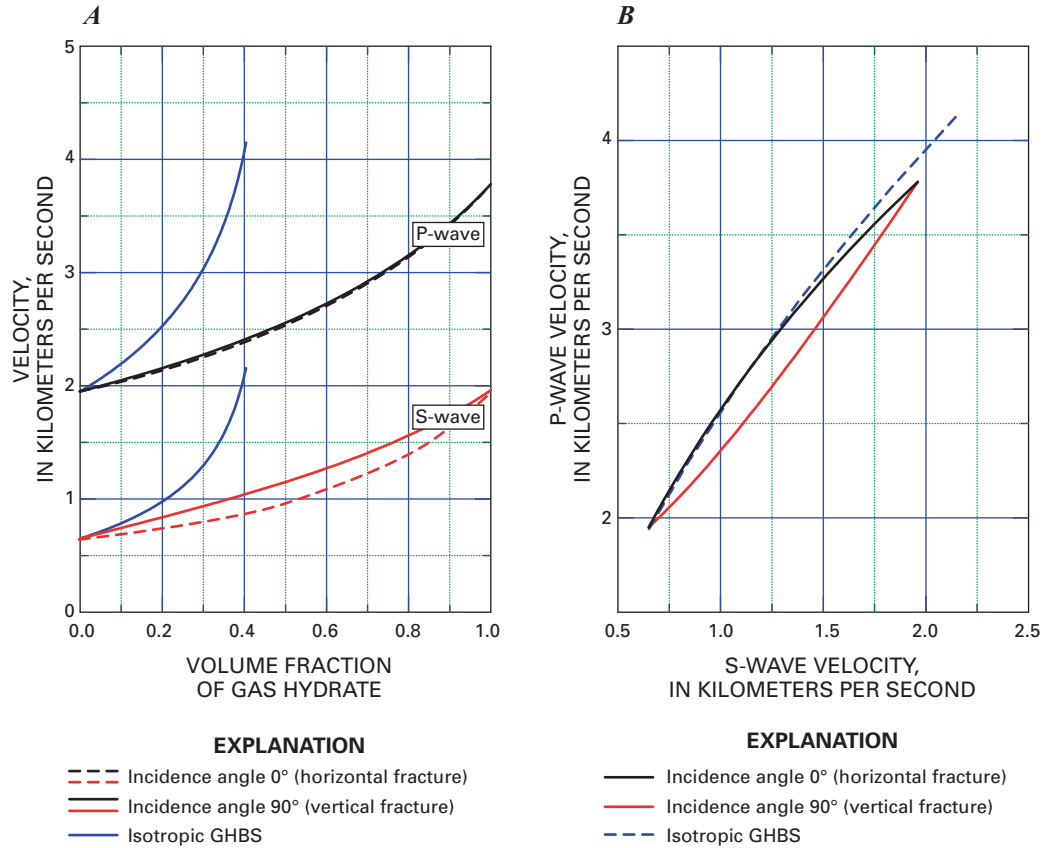


Figure 7. Modeling elastic velocities for anisotropic gas hydrate-bearing sediments (GHBS) caused by fracture in clean sandstone and for isotropic GHBS. *A*, Relation between gas hydrate saturation and bulk volume of gas hydrate in sediments. *B*, Relation between S-wave velocity and P-wave velocity.

of gas hydrate. Because the maximum volume fraction for the isotropic medium is limited by porosity, maximum volume fraction for the isotropic GHBS is 0.4. For a given gas hydrate content, the isotropic velocities are always larger than anisotropic velocities, and the orientation of fractures is insensitive to the amount of gas hydrate. For a given velocity, gas hydrate saturations estimated from the isotropic GHBS model are much less than those estimated from the anisotropic GHBS model.

Figure 7*B* shows the relation between P- and S-wave velocities of GHBS. The modeled relations using the isotropic and anisotropic GHBS are similar. Consequently, the relation between P- and S-wave velocities is not a good indicator for identifying fractured reservoirs in sandstone.

Because the fast and slow velocities of S-waves are almost identical, the presence of fractures in sandstone is difficult to detect on the basis of elastic velocities. The relation between velocity and saturation that is estimated independently from the velocity, such as saturations estimated from the nuclear magnetic resonance (NMR) log, can be used to identify the presence of fractures, although the orientation of the fractures may not be determined.

Gas Hydrate Saturations

The modeled velocities shown in figure 4 can be used to estimate gas hydrate saturations from measured velocities. Because the measured porosity ϕ_m is assumed to be an averaged value of the porosity of the water-saturated part of sediments and the porosity of the fracture filled with gas hydrate, the porosity of the water-saturated sediment is dependent on the amount of fracture (or gas hydrate). Therefore, an appropriate porosity for component 2 used to calculate velocities of water-saturated sediment is:

$$\phi_2 = \frac{\phi_m - V_h}{1 - V_h} \quad (10)$$

Equation 10 indicates that the porosity of water-saturated sediment is the same as the measured porosity of sediment in the absence of gas hydrate. As the volume of gas hydrate increases, the porosity of water-saturated sediment decreases to account for the presence of fractures.

The density porosity calculated from the density log assumes that the densities of gas hydrate and water are the same. Because the density of gas hydrate is less than that

of water, the calculated density porosity is not accurate for high-porosity sediments at high gas hydrate saturations. To compute accurate density porosity, the gas hydrate saturations are required. Consequently, estimating gas hydrate saturation and calculating density porosity are iterative processes, and the density porosity is updated using the estimated S_h as:

$$\phi_{new} = \phi_m \left[\frac{\rho_w - \rho_{ma}}{\rho_w(1 - S_h) + \rho_h S_h - \rho_{ma}} \right] \quad (11)$$

where ρ_w , ρ_h , and ρ_{ma} are density of water, gas hydrate, and grains, respectively. The updated ϕ_{new} is used for ϕ_m in equation 10.

Figure 8 shows wireline P- and S-wave velocities measured at drill site NGHP-10-10D along with logging-while-drilling (LWD) P-wave velocity at site NGHP-01-10A.

Wireline velocities were measured only for the gas hydrate interval. Figure 9 shows the gas hydrate saturations estimated assuming isotropic GHBS and anisotropic GHBS from wireline P- and S-wave velocities. The STPE parameters used for the estimation of gas hydrate assuming isotropic GHBS is $\alpha = 70(200 / d_f)^{1/3}$ and C_v calculated from the gamma log. The gas hydrate saturation estimated from the fracture model is an average value of fracture reservoir and water-saturated sediments and is calculated as $S_h = V_h / \phi_{new}$. As implied in figure 4, the gas hydrate saturations estimated assuming isotropic GHBS are between the two gas hydrate estimates: one assuming that the fracture is vertical and one assuming that the fracture is horizontal. Also shown as solid dots are the gas hydrate saturations estimated from the pressure cores. In this report, saturations estimated from the pressure cores are considered as reference saturations because pressure cores capture all the mass in samples. Gas hydrate saturations are estimated from the P-wave velocities by assuming vertical fractures agree better with those estimated from the pressure cores. However, the estimations from the pressure cores differ significantly from those estimated using the S-wave velocities assuming either isotropic or anisotropic GHBS. If the fractures are vertical, the orientation of the P-wave sensor does not matter for the velocity, because the particle motion of the P-wave is parallel to the fractures. On the other hand, the direction of particle motion of the S-wave is critical to estimate the velocities. Generally, it is not possible to accurately estimate gas hydrate saturation from the S-wave velocities without knowing the orientation of the S-wave sensor relative to the fracture. Because the orientation of the dipole sonic imager (DSI) tool is not available, it is assumed that the motion of the wireline S-wave is parallel to the fracture plane. That is, V_s^H shown in equation 4 is used as the theoretical S-wave velocity. Overall, the gas hydrate saturations assuming horizontal fracture agree well with those estimated assuming isotropic GHBS and agree well with that estimate from the pressure cores except the core at ≈ 74 mbsf.

Discussion

The model results indicate that the elastic characteristics of fractures in shale differ significantly from those in sandstone. The primary reason is that the elastic properties of pure methane are much higher than those of high-porosity shales, which are the most common host sediments for shallow marine GHBS, such as GHBS at drill site NGHP-01-10. On the other hand, the elastic velocities of sandstone reservoirs at moderate depths, such as GHBS at the Alaska North Slope, are much higher than those of shale at shallow depth; anisotropic effects of fractures in sandstone are not significant. However, to analyze the anisotropic elastic velocities due to fractures, the velocities of host sediments at a given locality should be accurately incorporated.

Gas hydrates at drill site NGHP-01-10 were inferred from LWD and wireline logs, and from recovered cores using X-ray images, cold spots, chlorinity anomaly, and moussey and soupy textures (Collett and others, 2008). The gas hydrate existed as solid nodules, high-angle and subhorizontal veins as fracture fill, and was disseminated throughout the cores. However, large differences exist among the various gas hydrate saturations. For example, the maximum gas hydrate saturation estimated from the pressure core is about 25.4 percent at

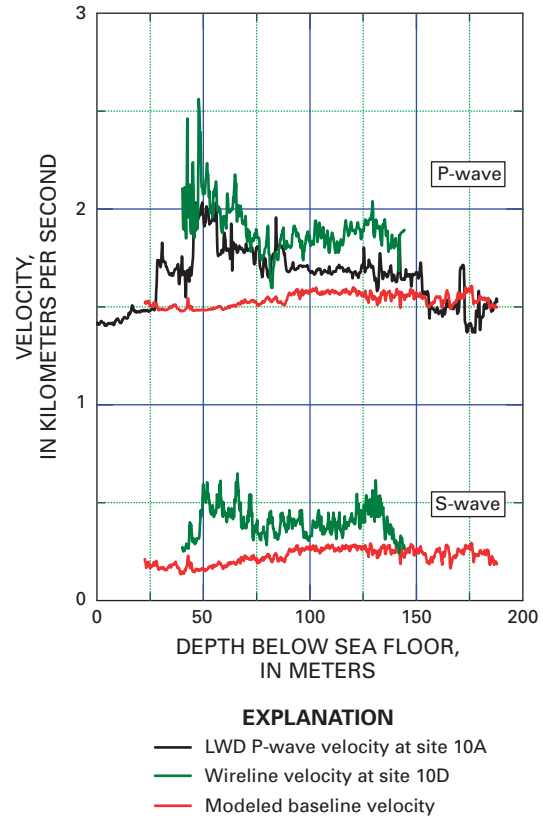


Figure 8. Measured logging-while-drilling (LWD) P-wave velocities at drill site NGHP-01-10A and wireline P- and S-wave velocities at drill site NGHP-01-10D. Calculated baseline velocities are shown in red.

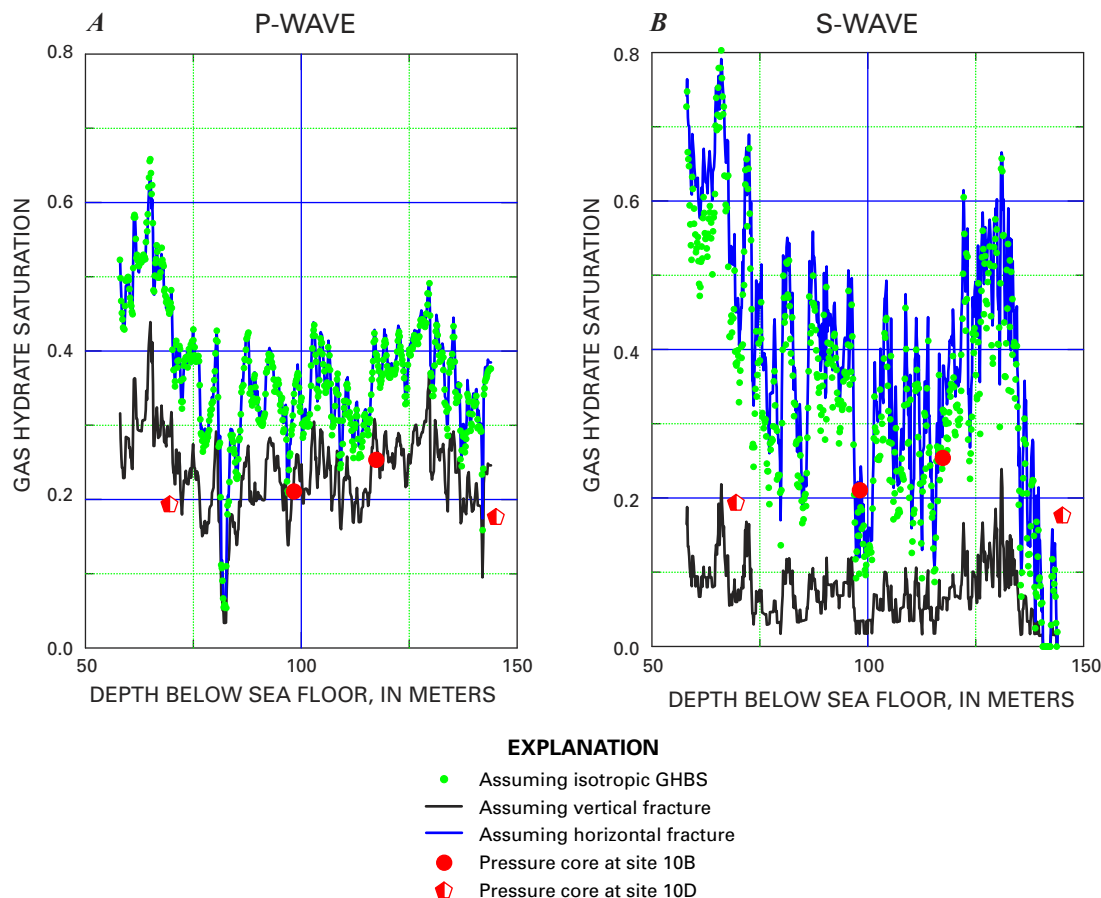


Figure 9. Gas hydrate saturations estimated from wireline velocities at drill site NGHP-01-10D assuming both anisotropic (vertical and horizontal fractures) and isotropic gas hydrate-bearing sediments (GHBS). *A*, From P-wave velocity. *B*, From S-wave velocity.

the depth of 117.4 mbsf at drill site NGHP-01-10B, which is shown in figure 9. The highest gas hydrate saturation assuming isotropic GHBS is about 80 percent, which was estimated from the LWD resistivity log as indicated in the initial report (Collett and others, 2008) or from S-wave velocities shown in figure 9.

Gas hydrate saturations estimated from the pressure cores depend only on the bulk volume of gas hydrate irrespective of whether gas hydrate fills the fractures or is dispersed throughout the sediments. The cumulative random measurement error in gas hydrate saturations from the pressure cores at drill site NGHP-01-10 is estimated slightly less than 5 percent (95-percent confidence interval) (M. Holland, written commun., 2009). The largest contributions of error are from porosity and ambient pressure. However, the loss of pressure and mass during core recovery occurred in only one sample at drill site NGHP-01-10 (Collett and others, 2008). Also, porosity was accurately measured. Therefore, these estimates can be used as reference saturations. Gas hydrate saturations from wireline P-wave velocities, assuming high-angle fractures, generally agree with those estimated from the pressure core, but are not consistent to those estimated from S-wave velocities due to uncertainties in the particle motion relative to the fracture orientation.

Conclusions

The anisotropic elastic velocities caused by fractures filled with 100-percent gas hydrate are modeled using laminated media and are applied to analyze velocities measured at drill site NGHP-01-10. From this study, the following conclusions can be drawn.

Fractures in shale at shallow depth filled with 100-percent gas hydrate significantly affect elastic velocities, particularly S-wave velocity, relative to the orientation of the fractures. On the other hand, velocities for fractured sandstone gas hydrate reservoirs are insensitive to the fracture orientation. However, to analyze the elastic velocities for fractured reservoirs, the velocities of host sediment at a given locality should be accurately modeled.

Assuming that velocities for gas hydrate-bearing sediments (GHBS) are isotropic could overestimate gas hydrate saturations for the fractured reservoirs, and the overestimation generally increases as the dip of the fracture increases.

An anisotropic GHBS model assuming two-component laminated media—one component is a fracture filled with 100-percent gas hydrate, and the other component is the isotropic water-saturated sediment—adequately predicts anisotropic velocities at drill site NGHP-01-10.

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